

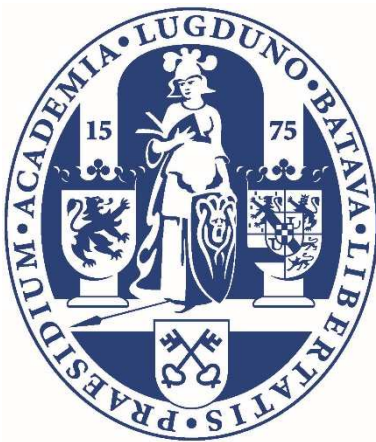
Quantum Optics: when do we need it?

Part 1: for optical field

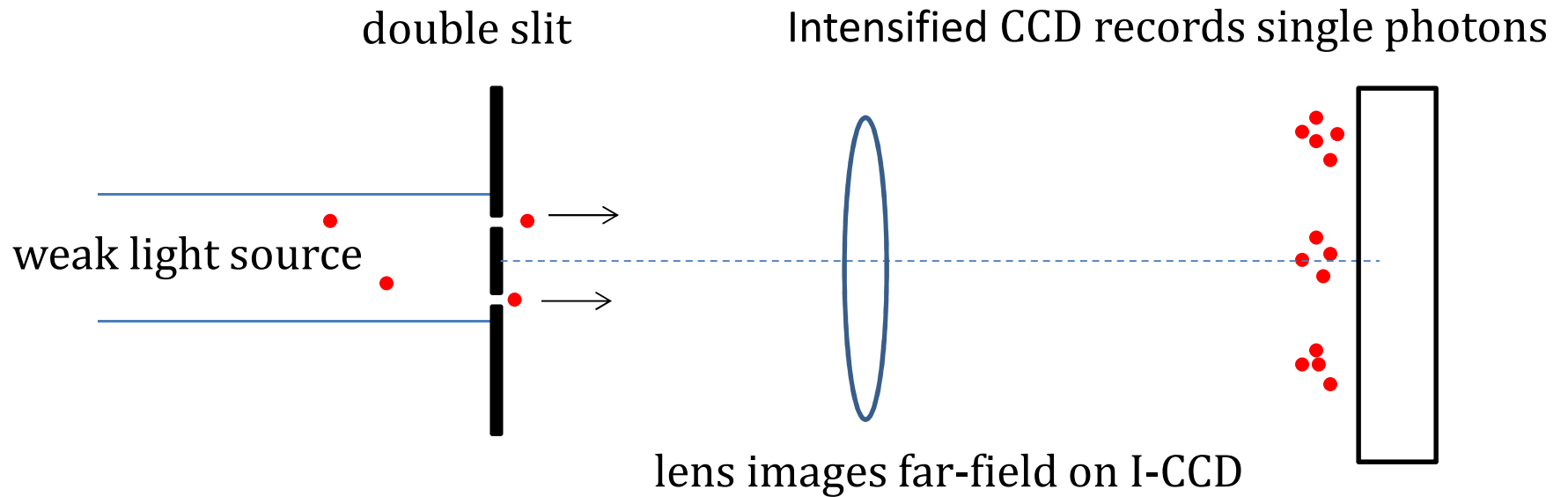
Martin van Exter

Huygens – Kamerlingh Onnes Laboratory

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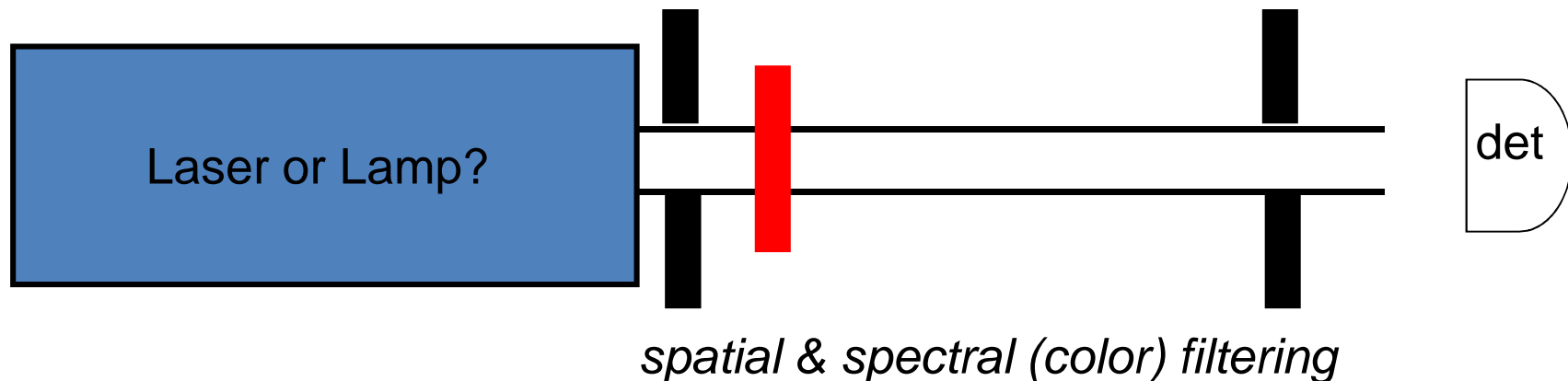


Example 1: Young's double slit with single photons



Experiment can be described without quantum optics

Example 2: How to distinguish laser from lamp light?



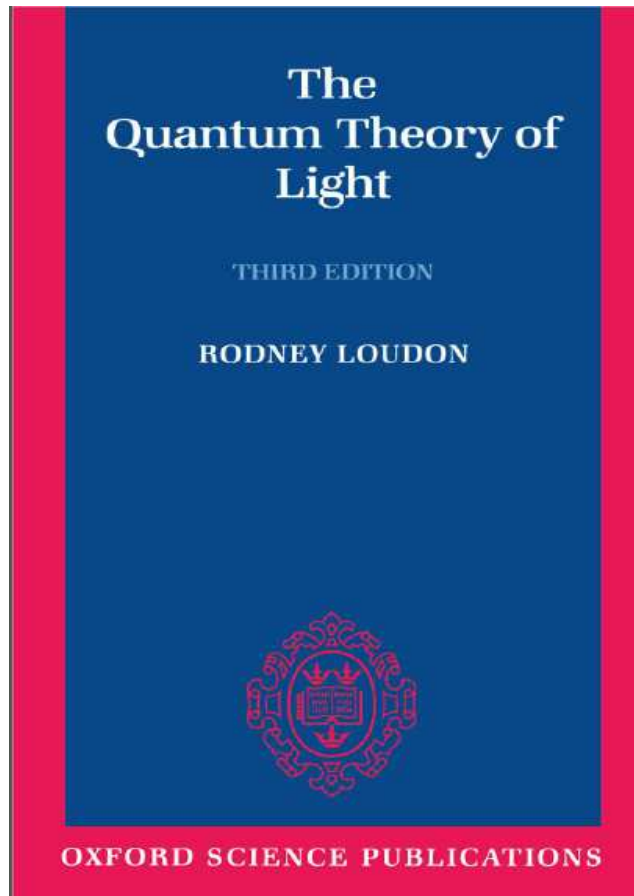
- “One-photon” optical properties (average intensity):
 - Temporal coherence & optical spectrum
 - Spatial coherence & intensity profile

Claim: “One-photon” properties don’t allow you to distinguish laser/lamp

- Correlation experiments between two detected photons are needed to distinguish different quantum states of light

‘Statistical Optics’ by J.W. Goodman

Single-mode vs. multi-mode quantum optics



Annihilation operator: \hat{a} or \hat{a}_i

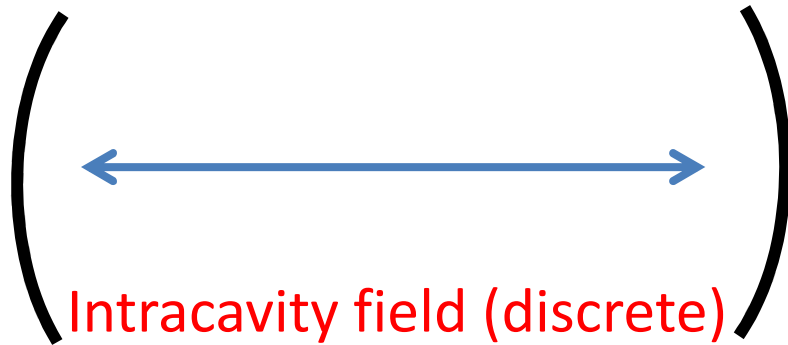
$$[\hat{a}_i, \hat{a}_j^\dagger] \equiv \hat{a}_i \hat{a}_j^\dagger - \hat{a}_j^\dagger \hat{a}_i = \delta_{i,j}$$

Annihilation operator: $\hat{a}(t)$ or $\hat{a}(\omega)$

$$(\hat{a}(\omega_1), \hat{a}(\omega_2)^\dagger) = \delta(\omega_1 - \omega_2)$$

- Single-mode = discrete mode in cavity (Chapter 5)
- Multi-mode = continuum of modes in free space (Ch. 6)

Textbooks often discuss discrete Q states of light



Intracavity field (discrete)

can be more quantum
than output field



Output field (continuous)

is affected by reflections
of vacuum field/fluctuations

Theory for “single-mode” quantum optics

- Single discrete mode: $[\hat{a}, \hat{a}^\dagger] = \hat{a}\hat{a}^\dagger - \hat{a}^\dagger\hat{a} = 1$

1. Number states:

$$|\psi\rangle = |n_0\rangle \Rightarrow P_n = \delta_{n,n_0}$$

2. Coherent states:

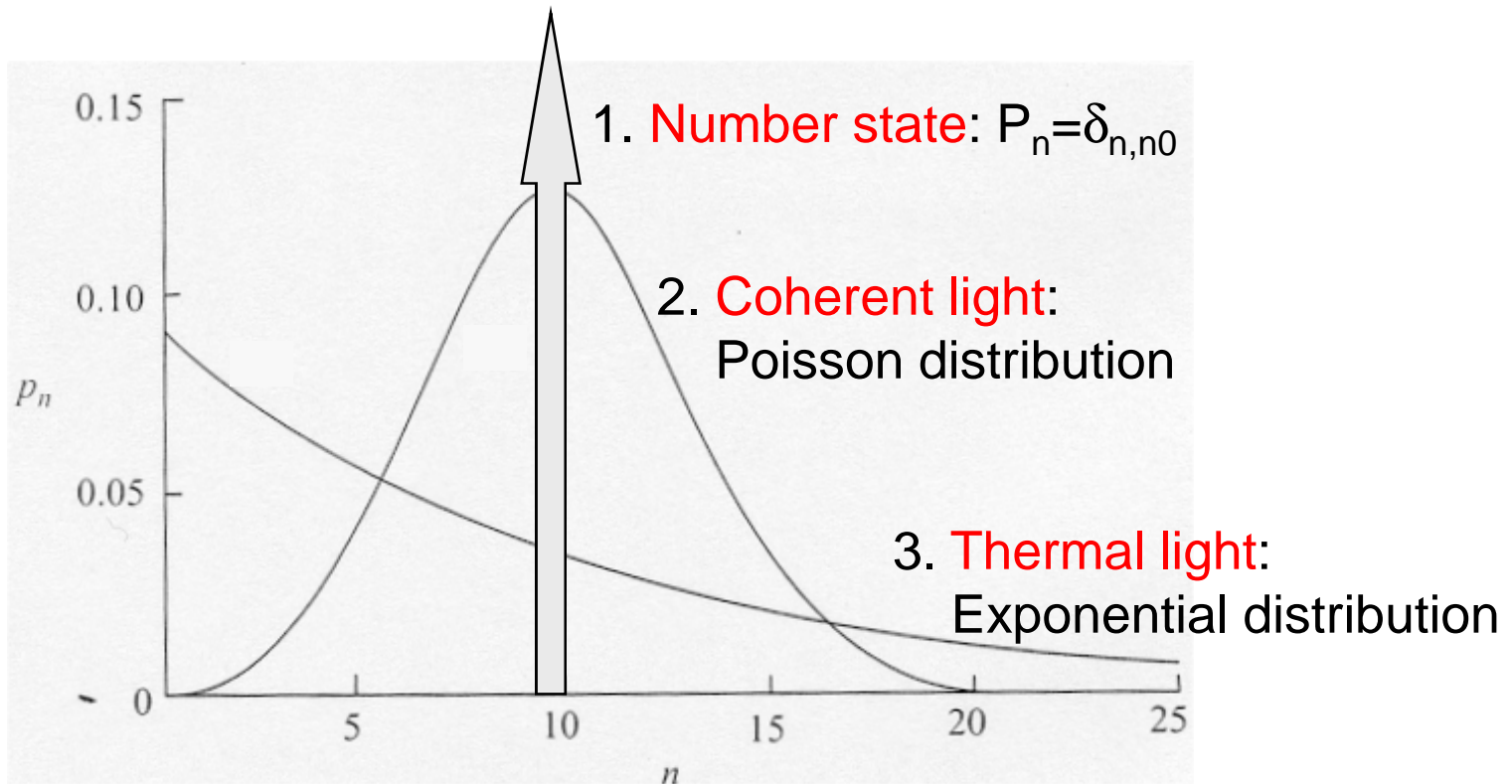
$$|\psi\rangle = |\alpha\rangle = e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle \Rightarrow P_n = e^{-|\alpha|^2} \frac{|\alpha|^{2n}}{n!} \text{ (Poisson distribution)}$$

3. Thermal light:

$$\rho = |\psi\rangle\langle\psi| = \sum_{n=0}^{\infty} P_n |n\rangle\langle n|$$

with $P_n \propto \exp\left(-n \frac{h\nu}{kT}\right)$ (Exponential distribution)

Three standard (single-mode) quantum states



“Multi-mode optics”

- Still single discrete transverse optical mode (fixed spatial structure)
- Continuous in frequency/time:

$$\left[\hat{a}(\omega), \hat{a}^\dagger(\omega') \right] = \delta(\omega - \omega')$$

- Temporal coherence:

$$g^{(1)}(\tau) = \frac{\langle E^-(t) E^+(t+\tau) \rangle}{\langle I(t) \rangle} \propto \langle \hat{a}^\dagger(t) \hat{a}(t+\tau) \rangle = \int I(\omega) \exp(-i\omega\tau) d\omega$$

- Intensity correlations:

$$g^{(2)}(\tau) = \frac{\langle :I(t) I(t+\tau): \rangle}{\langle I(t) \rangle^2} \propto \langle \hat{a}^\dagger(t) \hat{a}^\dagger(t+\tau) \hat{a}(t) \hat{a}(t+\tau) \rangle$$

- Only for thermal light (Gaussian statistics):

$$g^{(2)}(\tau) = 1 + |g^{(1)}(\tau)|^2$$

When do you need quantum optics?

- **Single-mode quantum optics**
is rarely relevant in free-space optics experiments
- **Multi-mode quantum optics** needs field quantization for:
 1. Direct observation of intensity noise (sensitive to loss)
 2. Photon-photon correlations (Hanbury Brown & Twiss)
 - 2b. Interference of two single photons (Hong, Ou & Mandel)

1. Semi-classical theory of photon detection

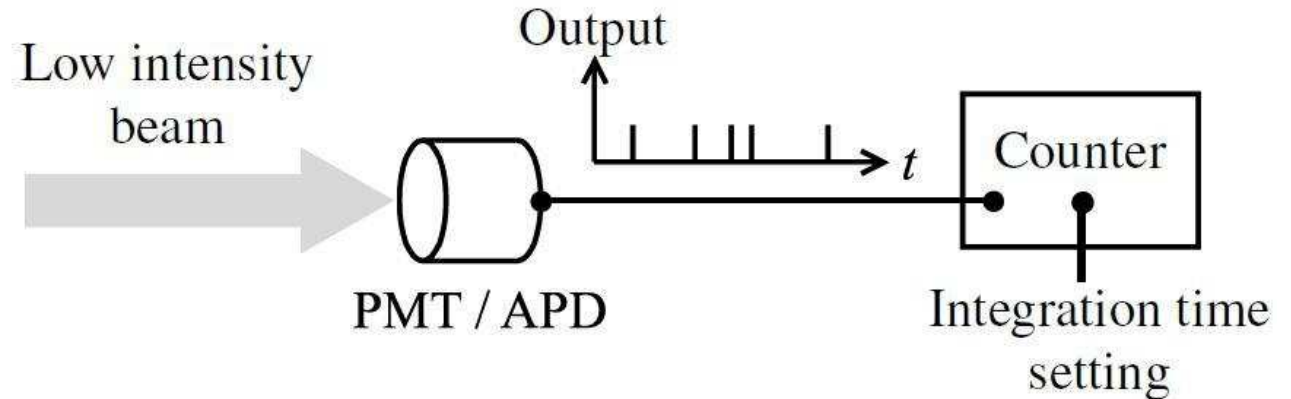
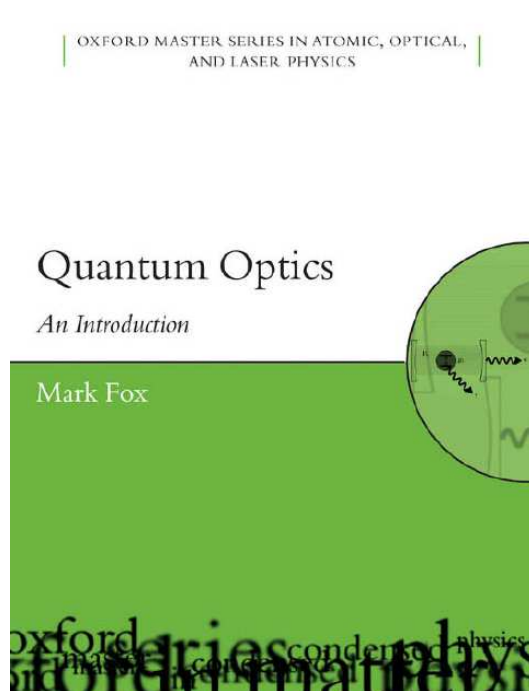
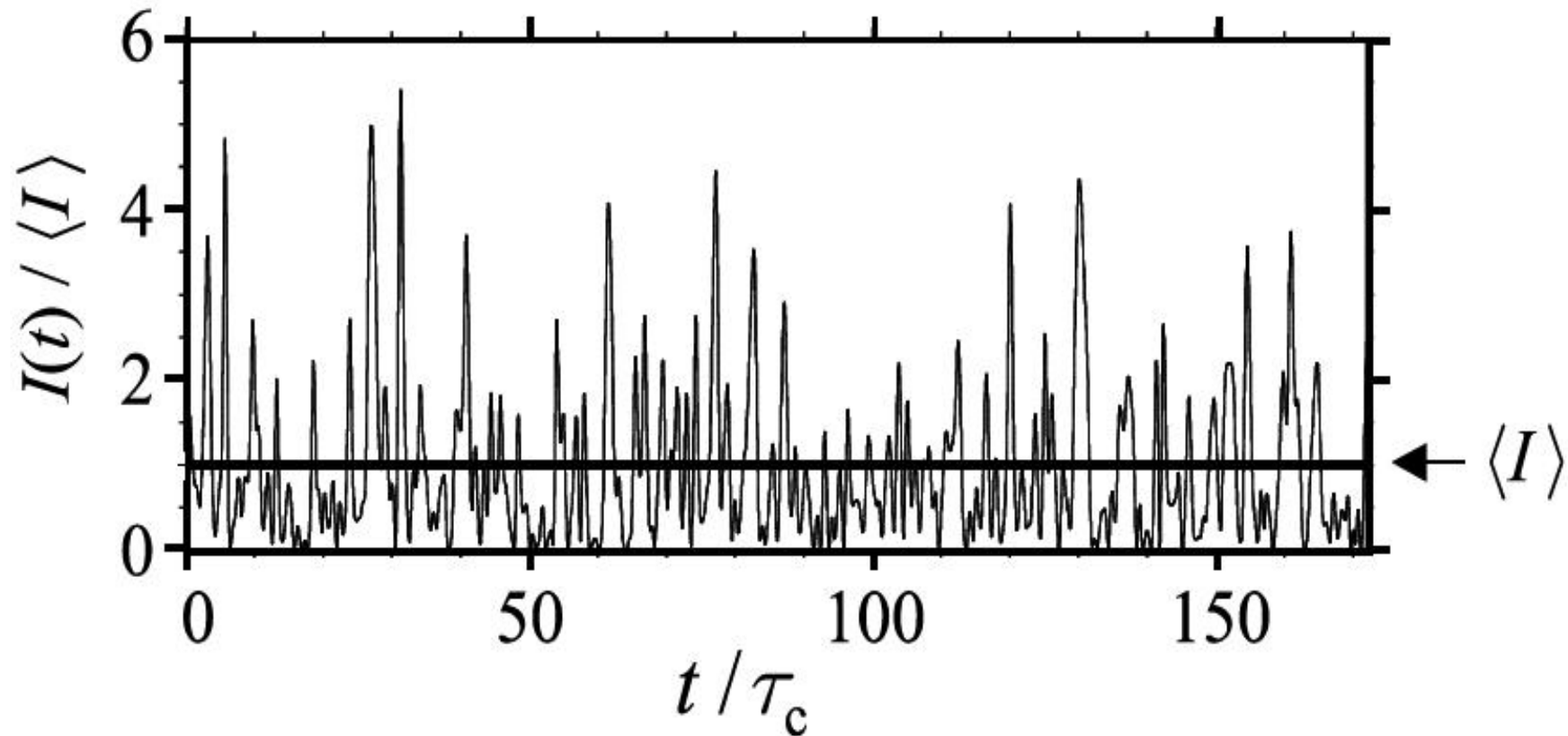


Fig. 5.1 of 'Quantum Optics' by M. Fox

- Semi-classical = treat field classically & only quantize detector

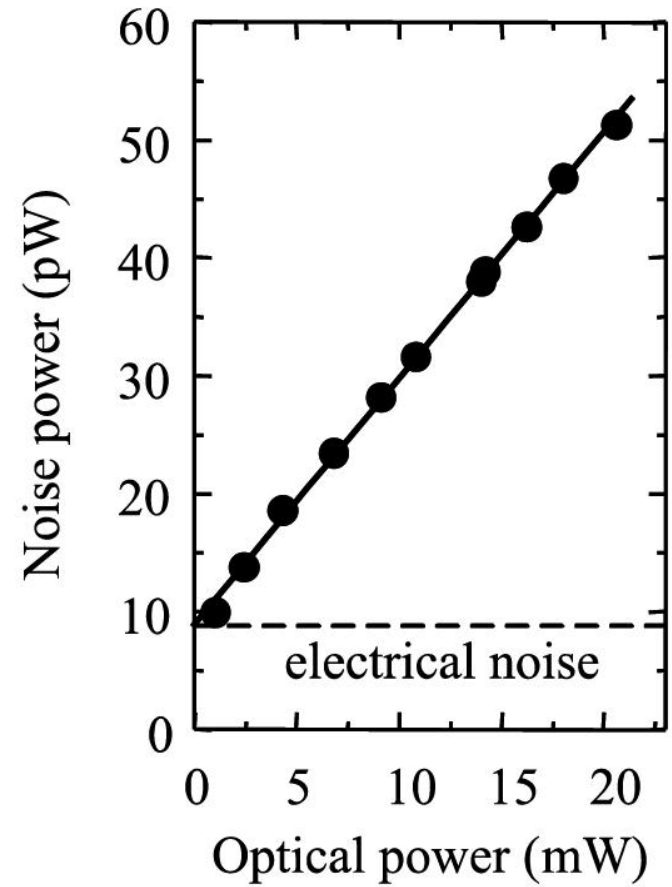
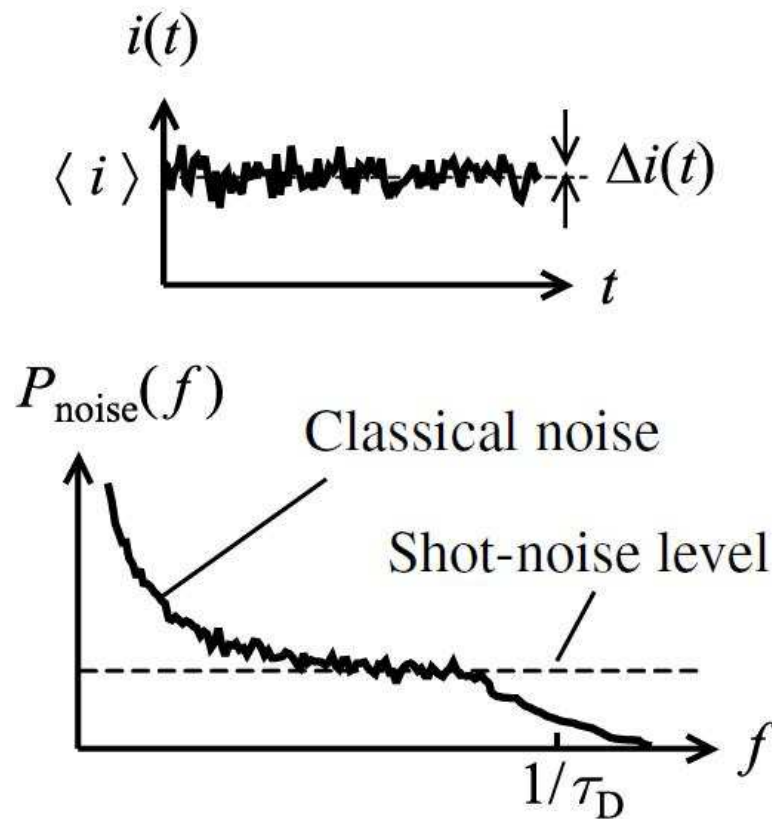
$$P(t; t + dt) = \eta \cdot \frac{P(t)}{h\nu} \cdot dt$$

Classical description of intensity fluctuations



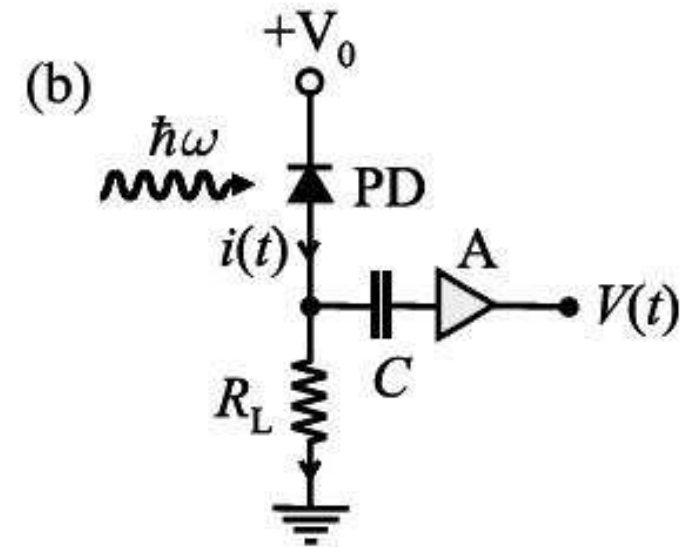
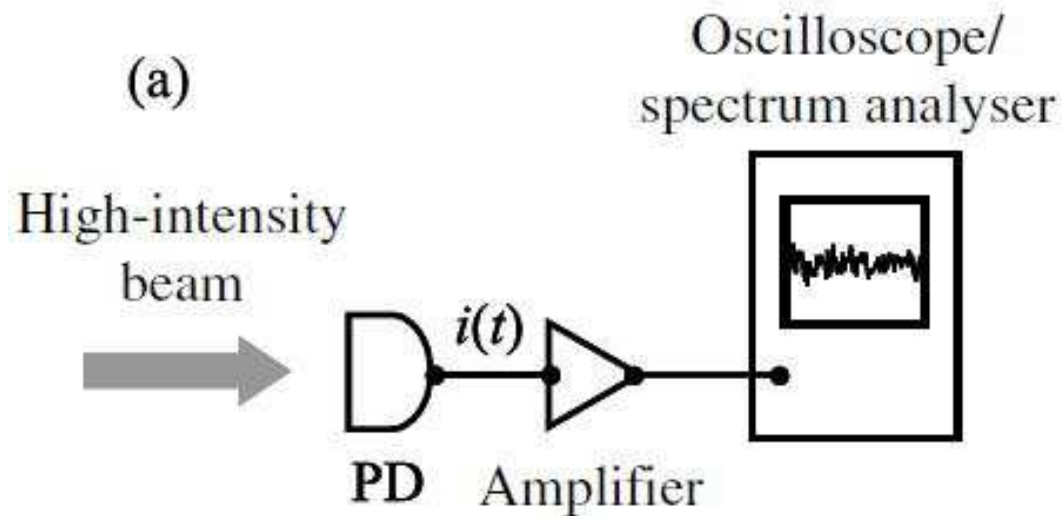
Thermal light exhibits **fluctuations of both optical phase and intensity**

Shot noise in the detection of light



Intensity noise $I_{rms} \propto \sqrt{I_{DC}}$

Shot noise can only dominate when $V_{det} > 50 \text{ mV}$



$$V_{electric\ noise} = \sqrt{4kTR_L \cdot \Delta f}$$

($\approx 0.9 \text{ nV}/\sqrt{\text{Hz}}$ in 50Ω at roomtemperature)

$$V_{shot\ noise} = R_L \cdot \sqrt{2qi \cdot \Delta f} \quad (\text{scales as } \sqrt{N})$$

$$V_{shot\ noise} > V_{electric\ noise} \Rightarrow V_{det} = R_L i > 2kT/q \quad 50 \text{ mV @ } 300K$$

Sub-Poissonian light

Franck-Hertz experiment: sub-shot noise intensity fluctuations

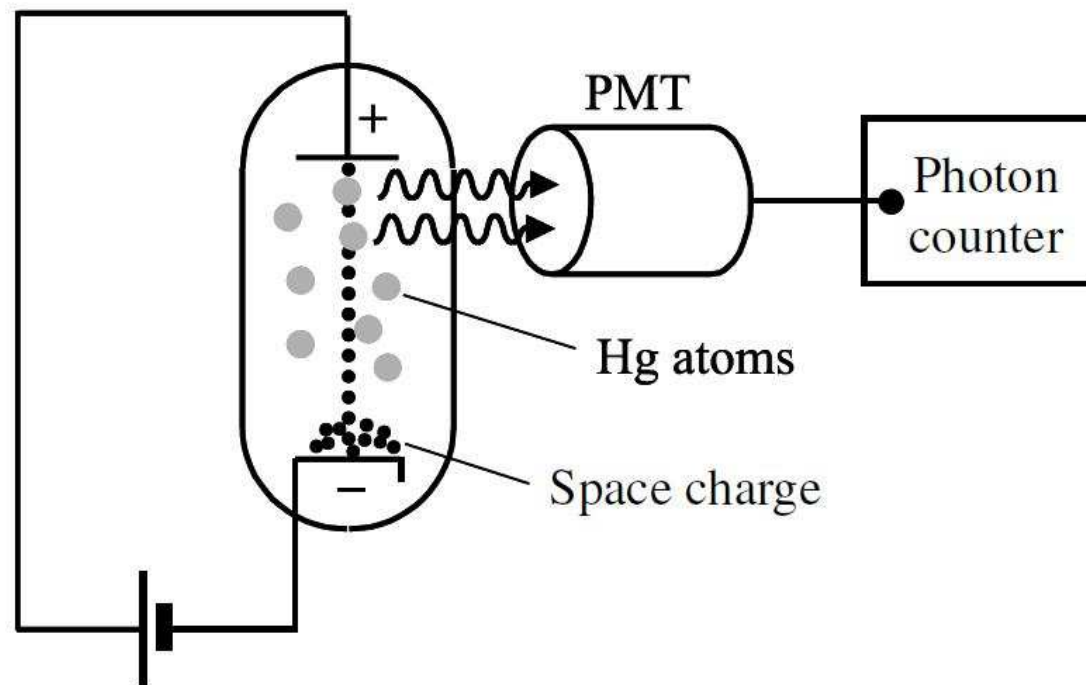
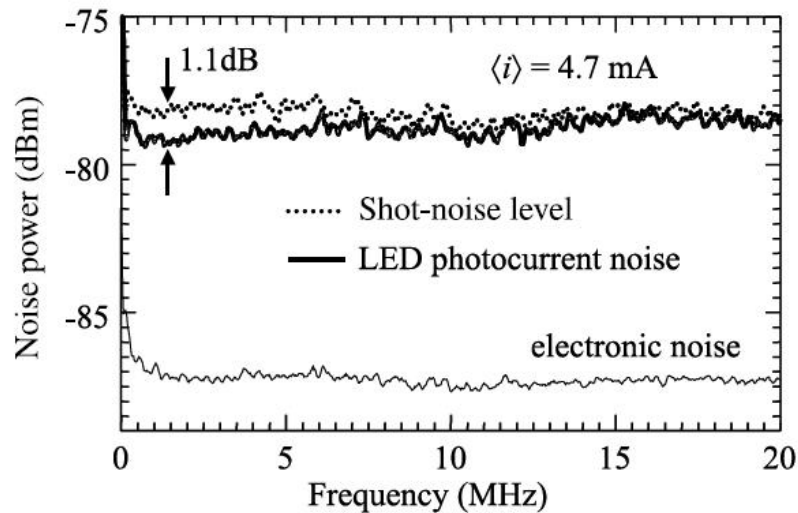
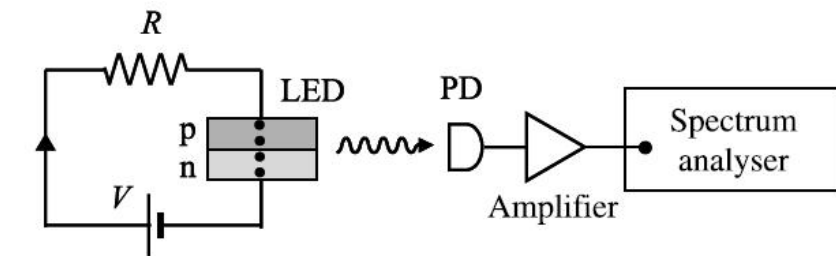


Fig. 5.14 of 'Quantum Optics' by M. Fox

Sub-Poissonian light

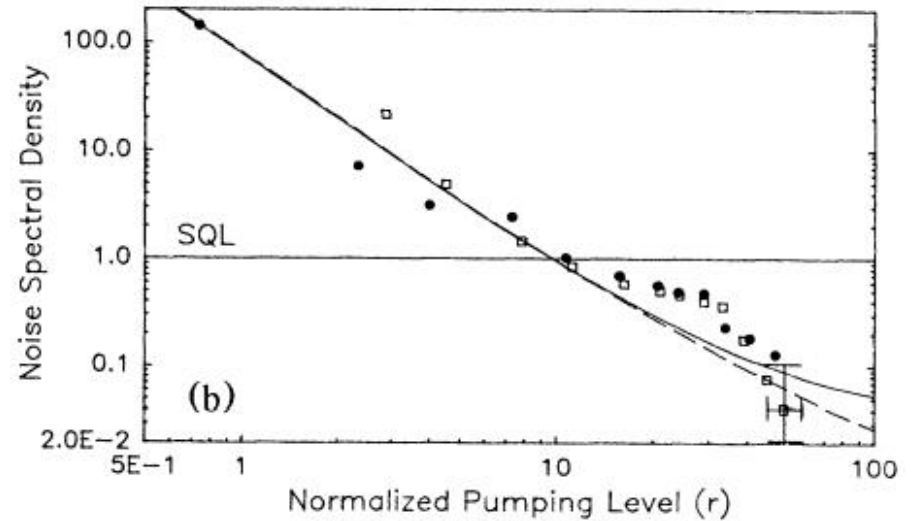
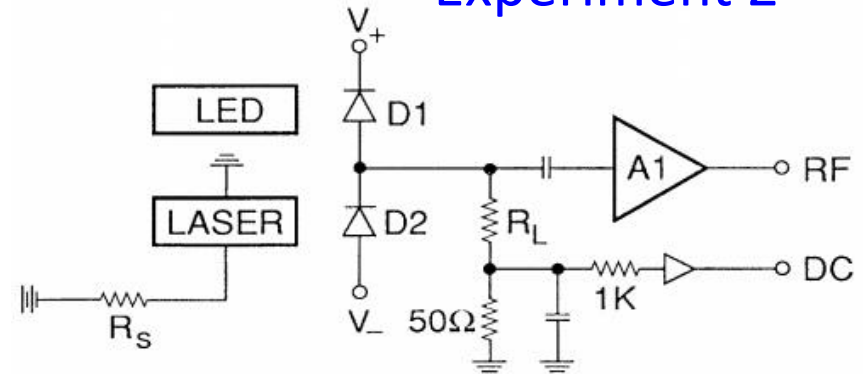
Electric current doesn't suffer from shot-noise

Experiment 1



M. Fox, *Quantum Optics*, Fig. 5.15

Experiment 2



W.H. Richardson et al., *PRL* 66,2867 (1992)

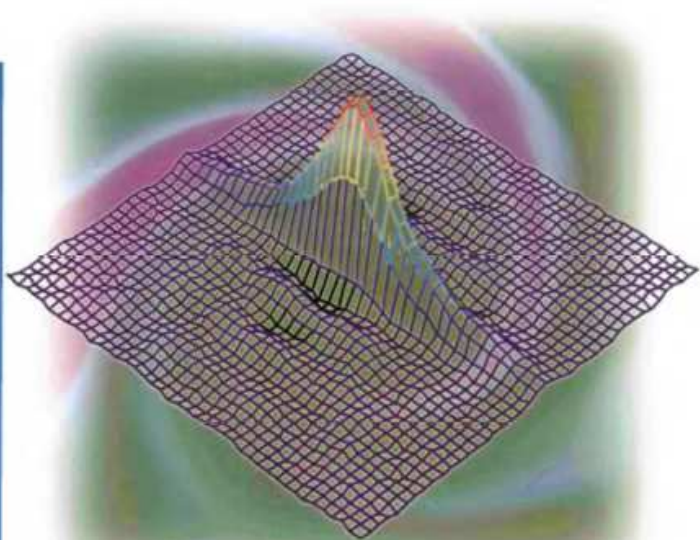
Optical loss kills the sub-Poissonian character

PHYSICS TEXTBOOK

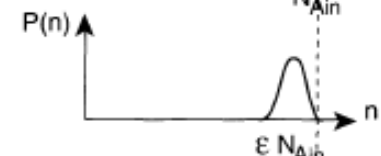
Hans-A. Bacher and Timothy C. Ralph WILEY-VCH

A Guide to Experiments in Quantum Optics

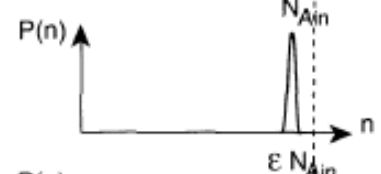
Second, Revised and Enlarged Edition



(a) Coherent or Poissonian input

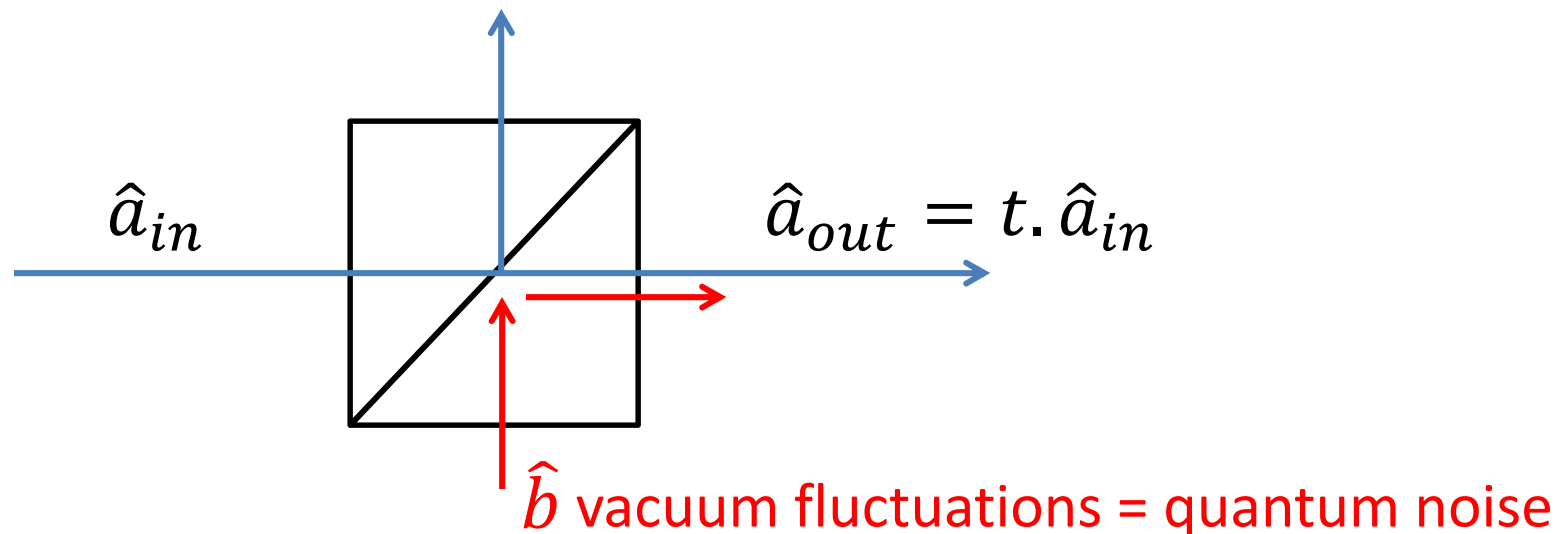


(b) Sub Poissonian input



Optical loss kills the sub-noise character

Optical loss introduces quantum noise, as it acts like beamsplitter



$$[\hat{a}_{out}, \hat{a}_{out}^\dagger] = |t|^2 \cdot [\hat{a}_{in}, \hat{a}_{in}^\dagger] = |t|^2 ??$$

$$\hat{a}_{out} = t \cdot \hat{a}_{in} + r \cdot \hat{b} \quad \text{with} \quad [\hat{b}, \hat{b}^\dagger] = 1$$

$$[\hat{a}_{out}, \hat{a}_{out}^\dagger] = |t|^2 + |r|^2 = 1$$

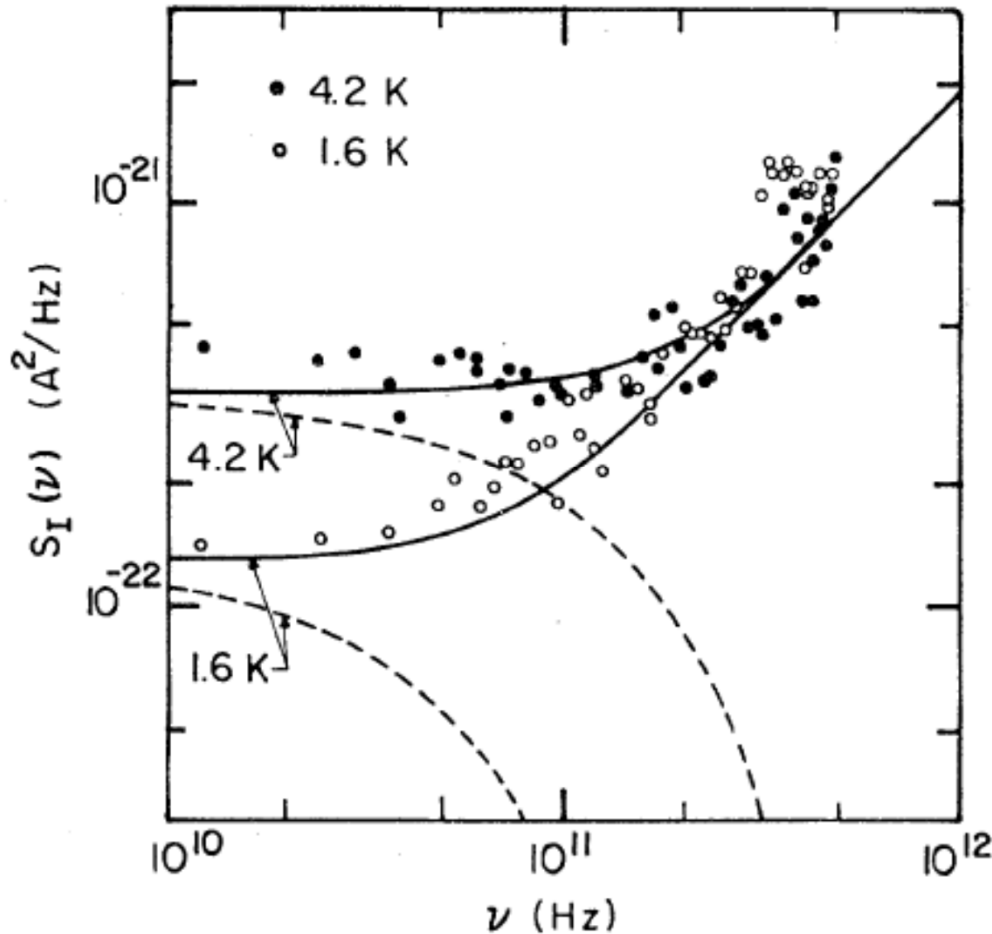
Quantum noise = Vacuum fluctuations

- Strength of the vacuum fluctuations depends on measurement:

$$\langle 0 | \hat{a}^\dagger \hat{a} | 0 \rangle = 0 \quad \langle 0 | \frac{1}{2} (\hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger) | 0 \rangle = \frac{1}{2} \quad \langle 0 | \hat{a} \hat{a}^\dagger | 0 \rangle = 1$$

- Vacuum fluctuations (in single transverse mode) :
 - 1 photon/second / per unit spectral bandwidth = 0.12 μ W /nm @ 800 nm

Quantum noise can be measured



Experiment: current noise in Josephson tunnel junction, mixed down to detection frequencies $(\hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger)$ – form

$$S_I(\nu) = \frac{4h\nu}{R} \left(\frac{1}{\exp(h\nu/kT) - 1} + \frac{1}{2} \right)$$

Limit $kT \gg h\nu$

$$S_I(\nu) = \frac{4kT}{R}$$

Limit $kT \ll h\nu$

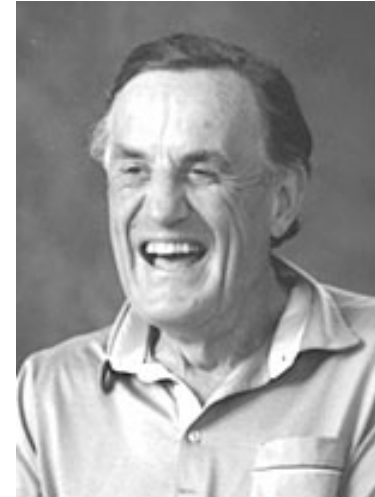
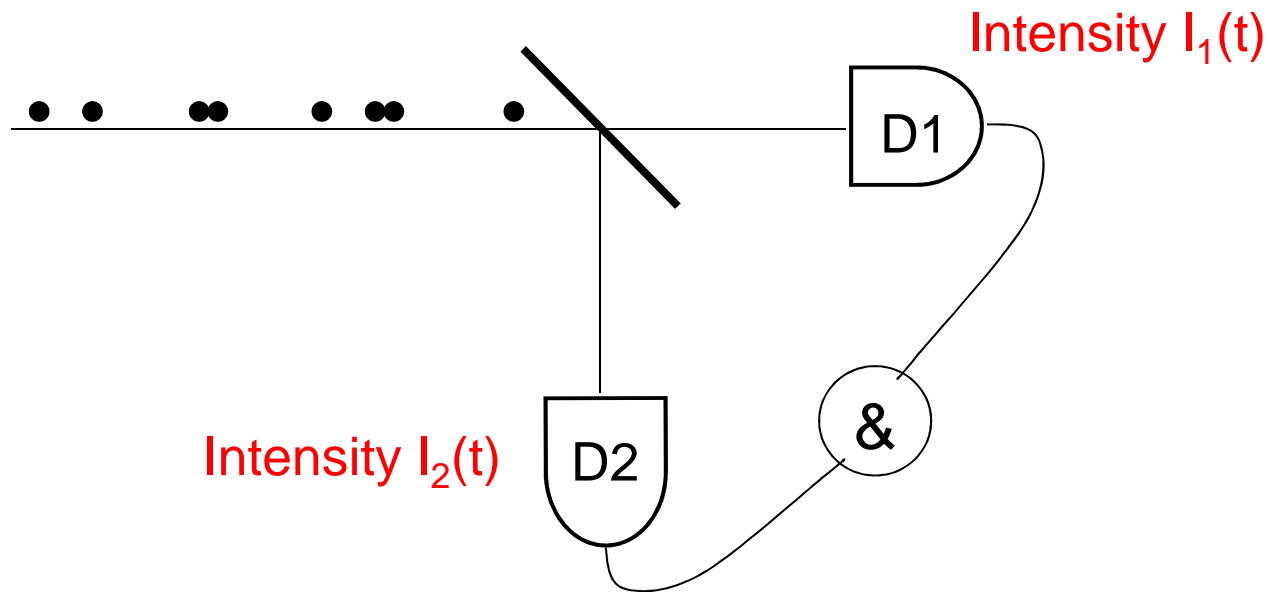
1/2 noise quantum per “time/frequency mode”

R.H. Koch et al., PRB 26, 74 (1982)

Introduction in book C.W. Gardiner, ‘Quantum Noise’

2. Intensity fluctuations & photon correlations

Hanbury Brown & Twiss experiment



Hanbury Brown

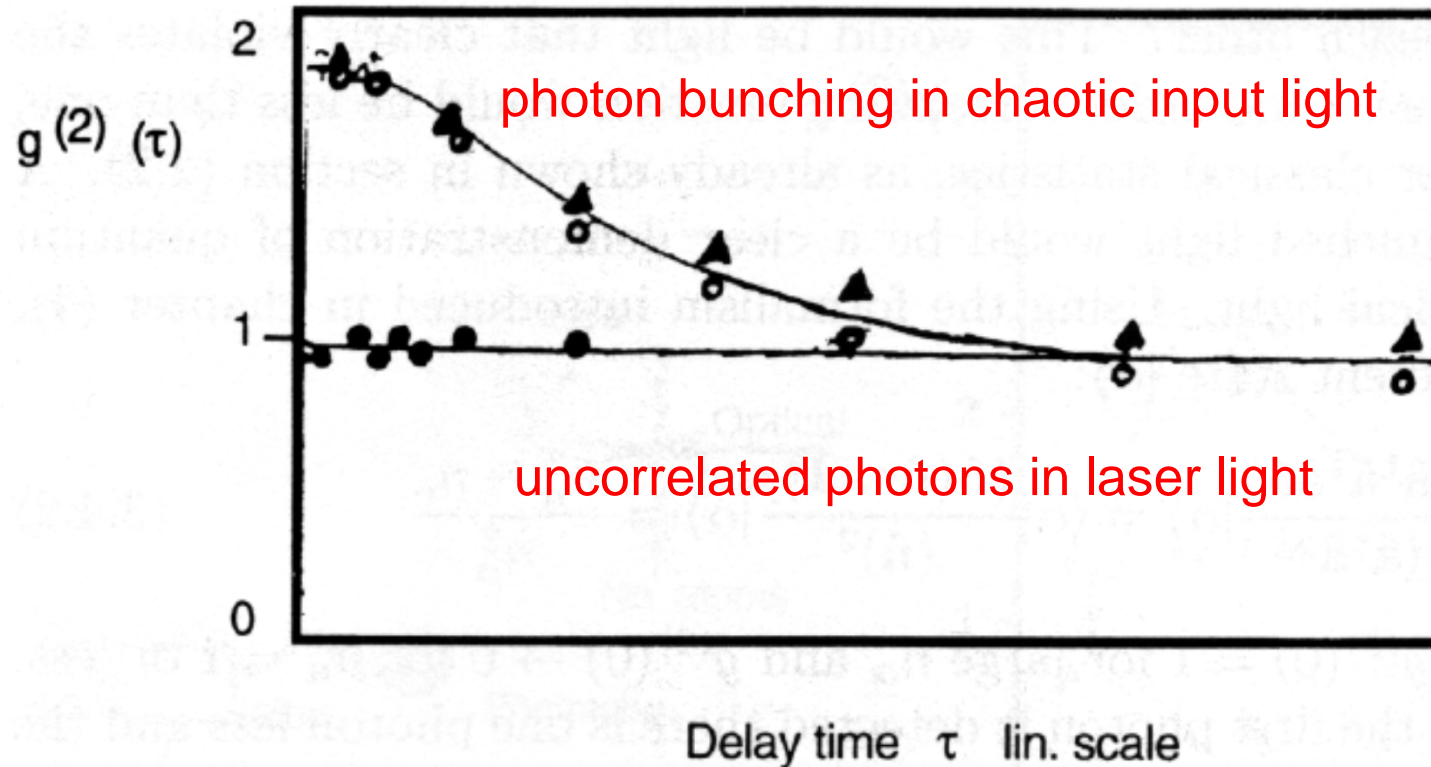


Twiss

Second-order (intensity-intensity) correlation function

$$g^{(2)}(\tau) \equiv \frac{\langle I(t+\tau)I(t) \rangle}{\langle I(t) \rangle^2}$$

Intensity fluctuations distinguish laser from chaotic light

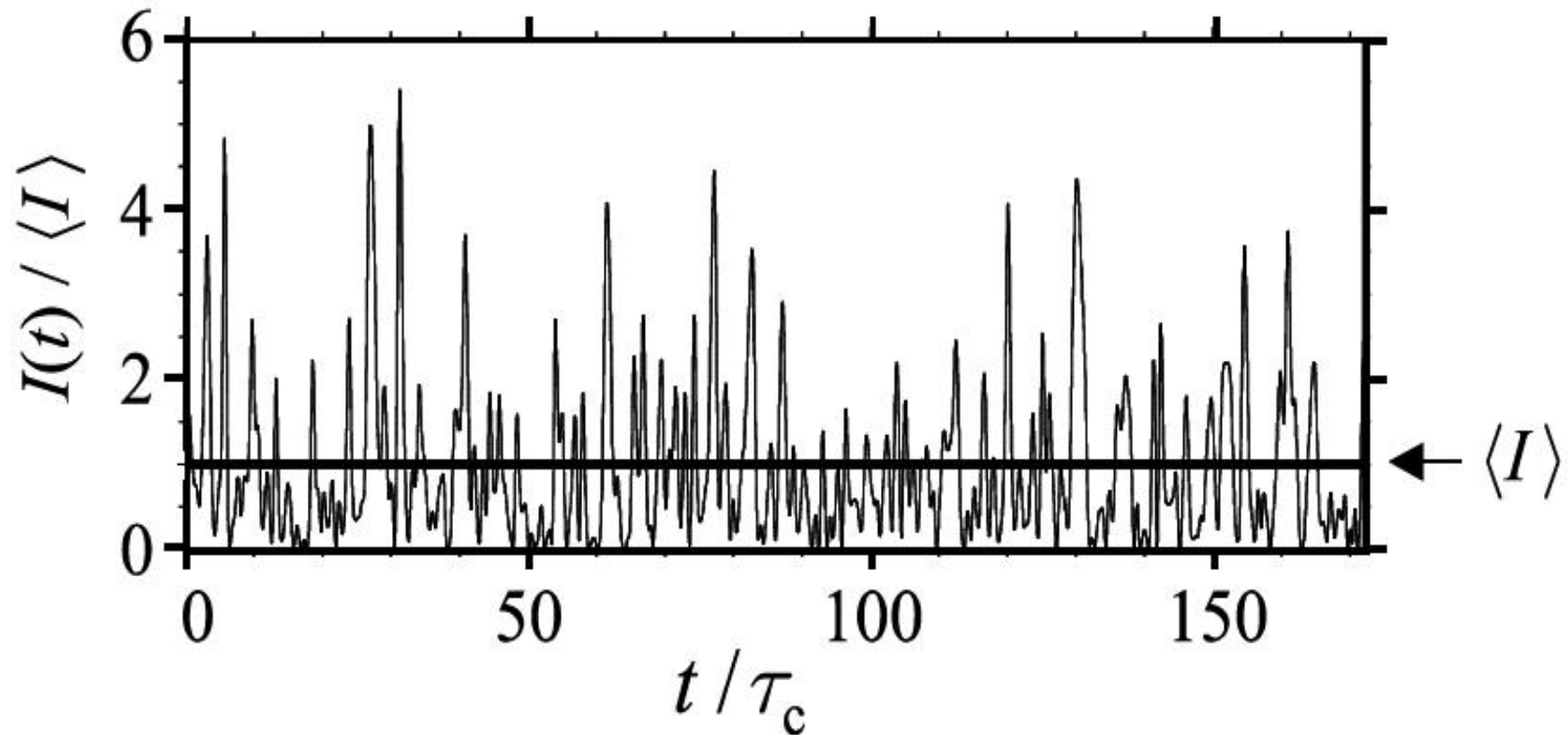


- Solid dots = laser above threshold
- other symbols = chaotic light

F.T. Arecchi et al., PRL 17, 260 (1966)

$$g^{(2)}(\tau) \equiv \frac{\langle I(t+\tau)I(t) \rangle}{\langle I(t) \rangle^2}$$

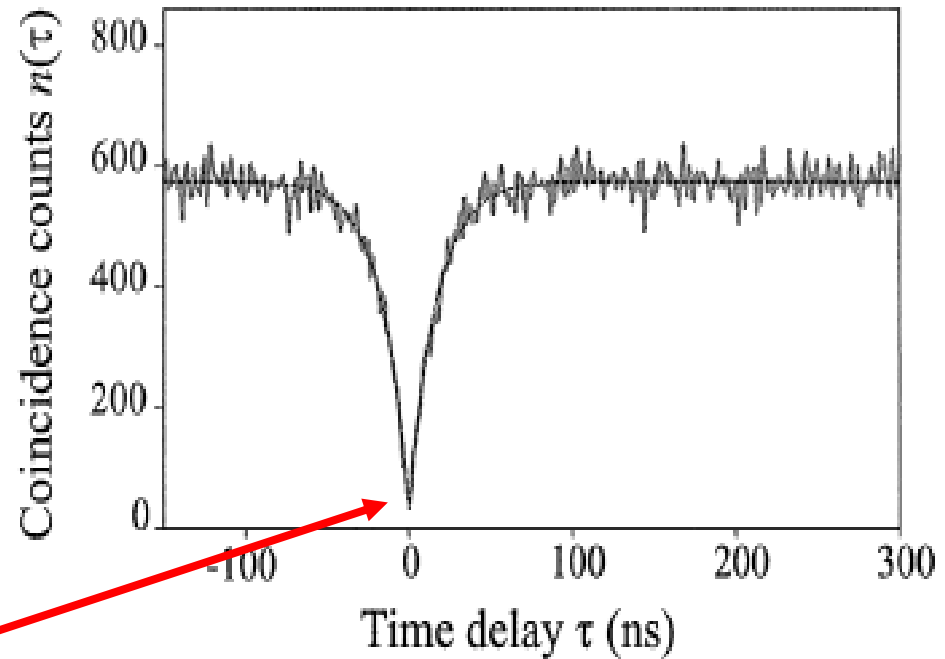
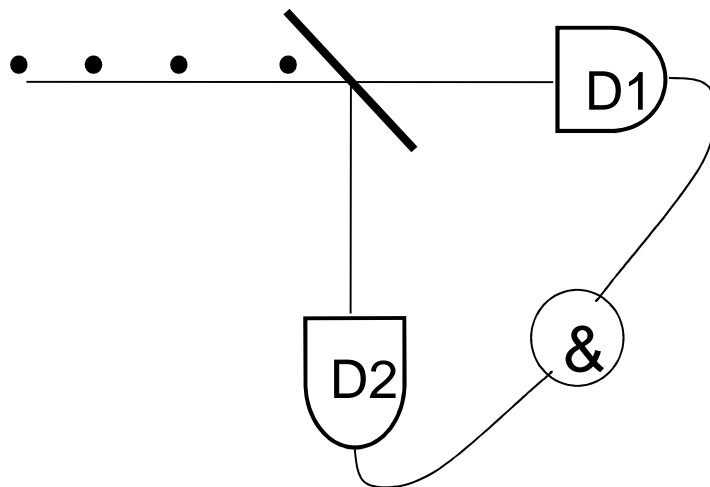
Classical description of intensity fluctuations



Thermal light exhibits **fluctuations of both optical phase and intensity**

Anti-bunching in fluorescence single CdSe quantum dot

Anti-bunching observed for light from single-photon sources



Anti-bunching at input: never two photons together!

G. Messin et al., Opt. Lett. 26, 1891 (2001)

Time-series of photon detection events



(a)

photon bunching:

$$P_{\text{conditional}}(t + \tau | t) \propto \bar{I} (1 + \exp(-\tau/\tau_c))$$



(b)

uncorrelated photons:

$$P_{\text{conditional}}(t + \tau | t) \propto \bar{I}$$

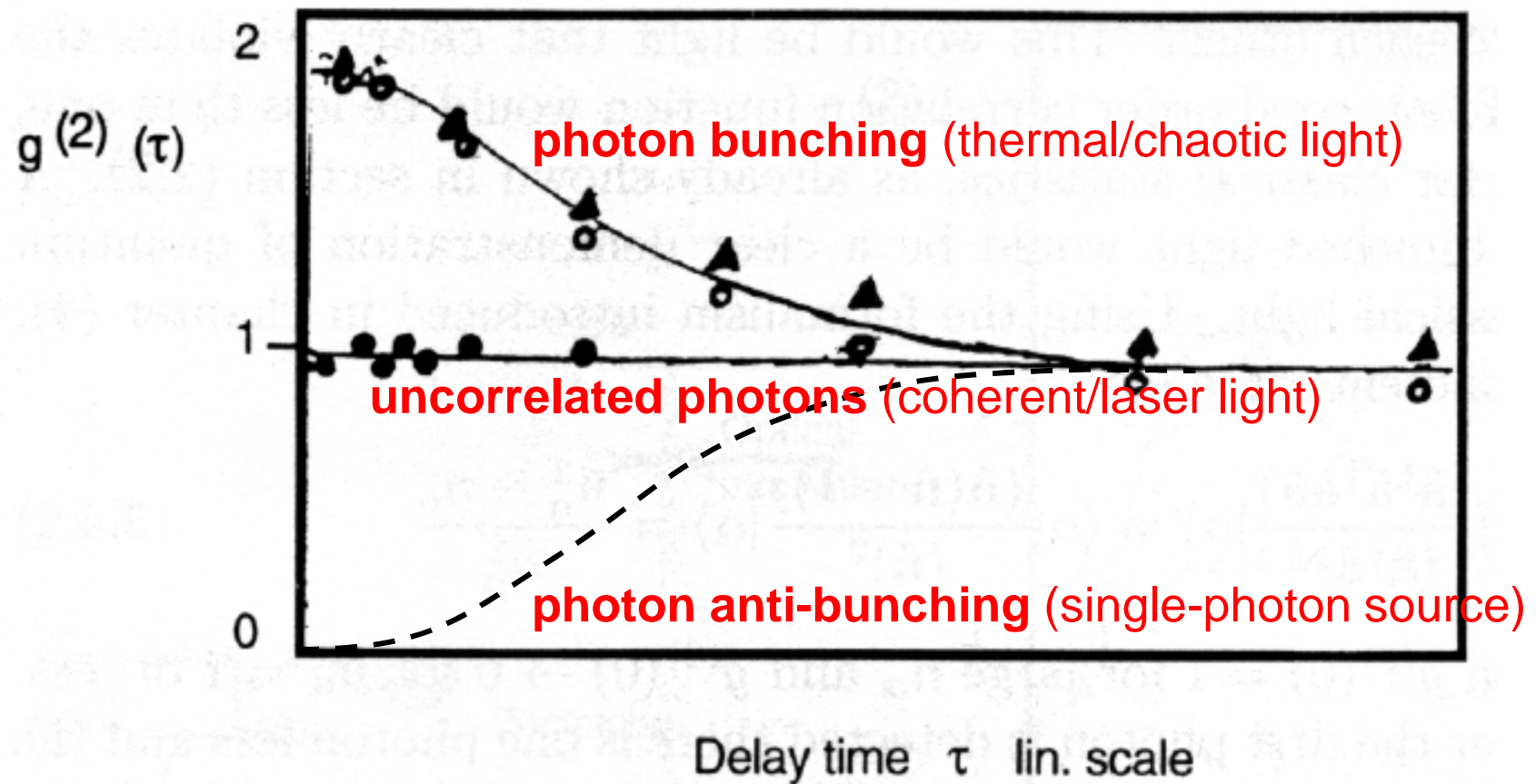


(c)

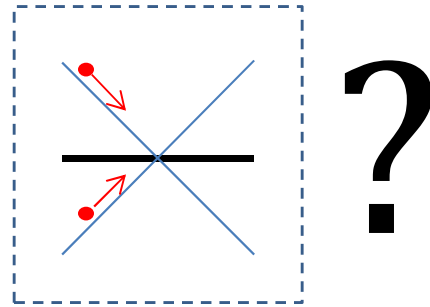
photon anti-bunching:

$$P_{\text{conditional}}(t + \tau | t) \propto \bar{I} (1 - \exp(-\tau/\tau_c))$$

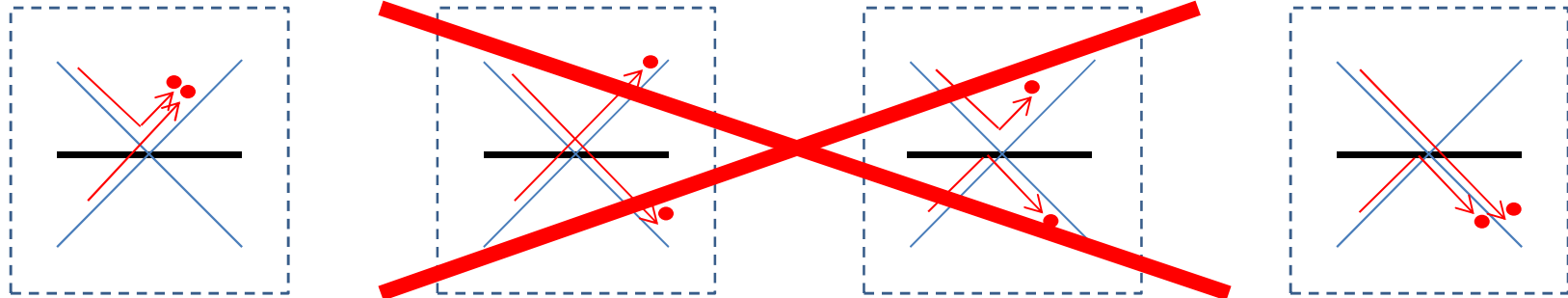
Intensity correlations for three quantum states of light



3. Interference between two identical single photons

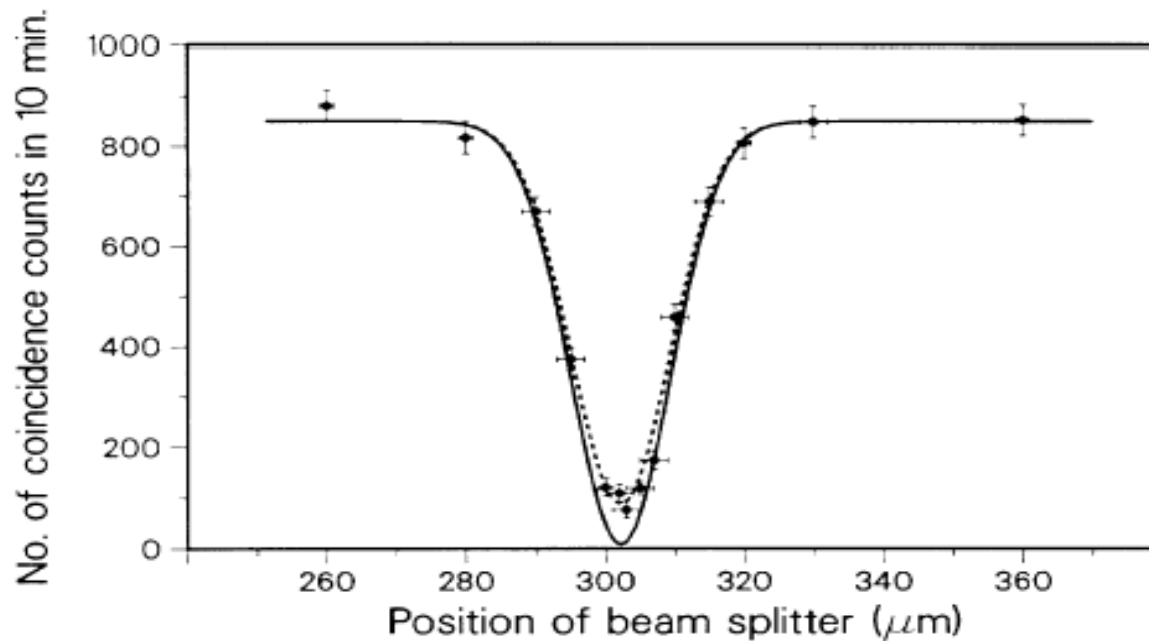
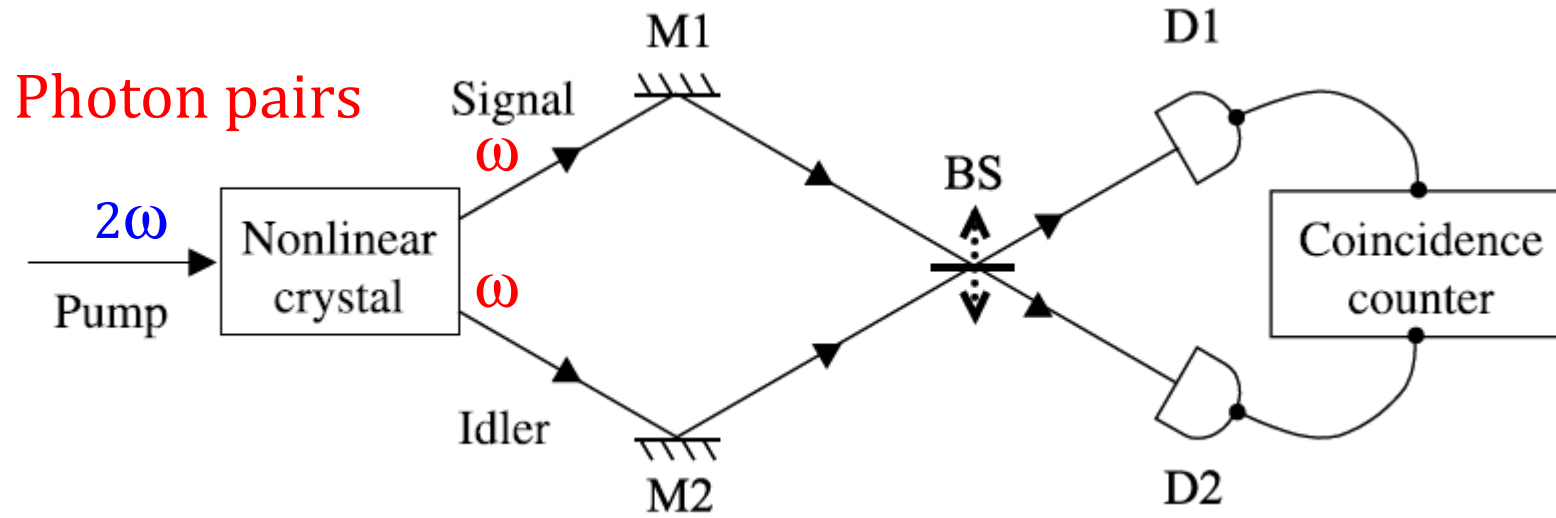


Four possible scenarios



Quantum prediction: no coincidences for identical photons

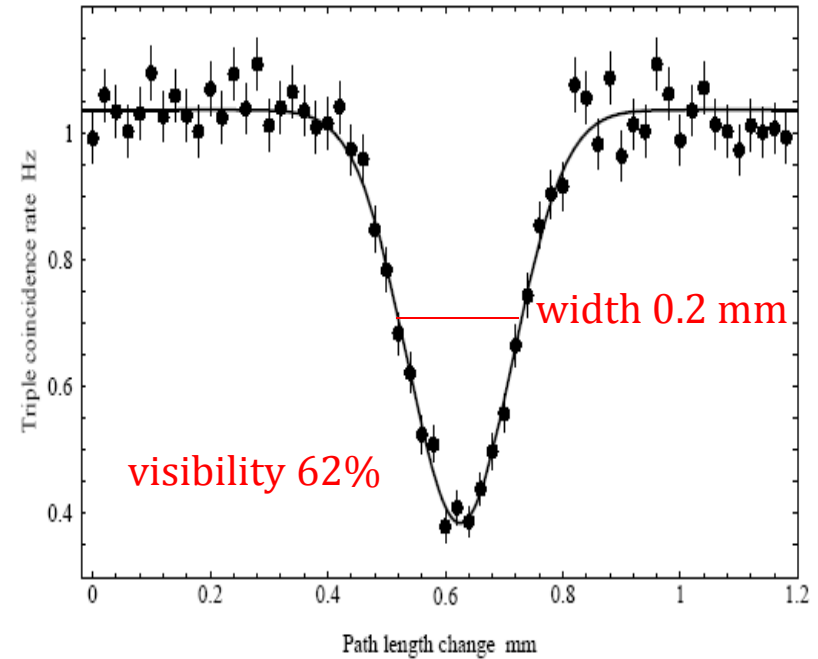
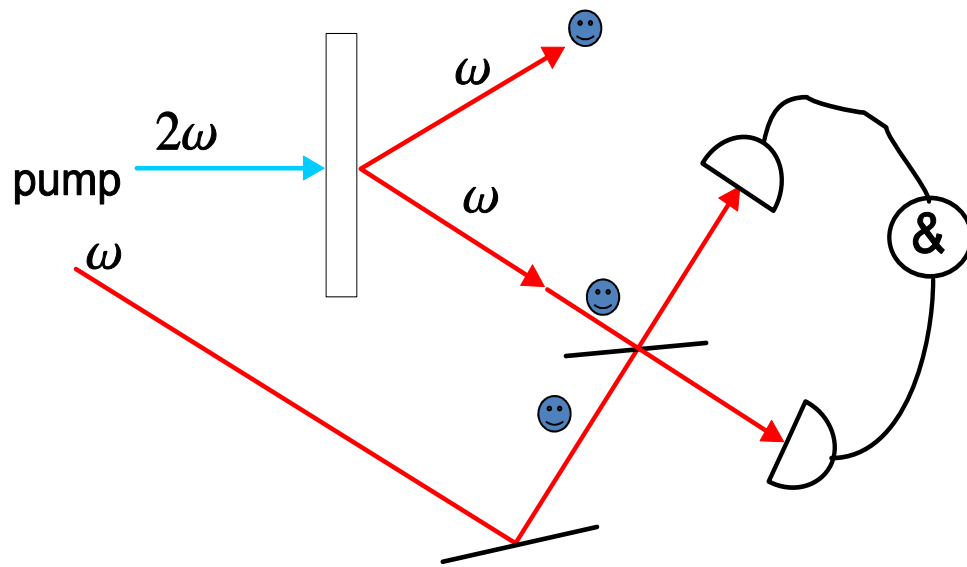
Photon bunching in Hong-Ou Mandel experiment



No fringes!

*Hong, Ou, Mandel,
PRL59, 2044 (1987)*

Interference between two 'independent' photons

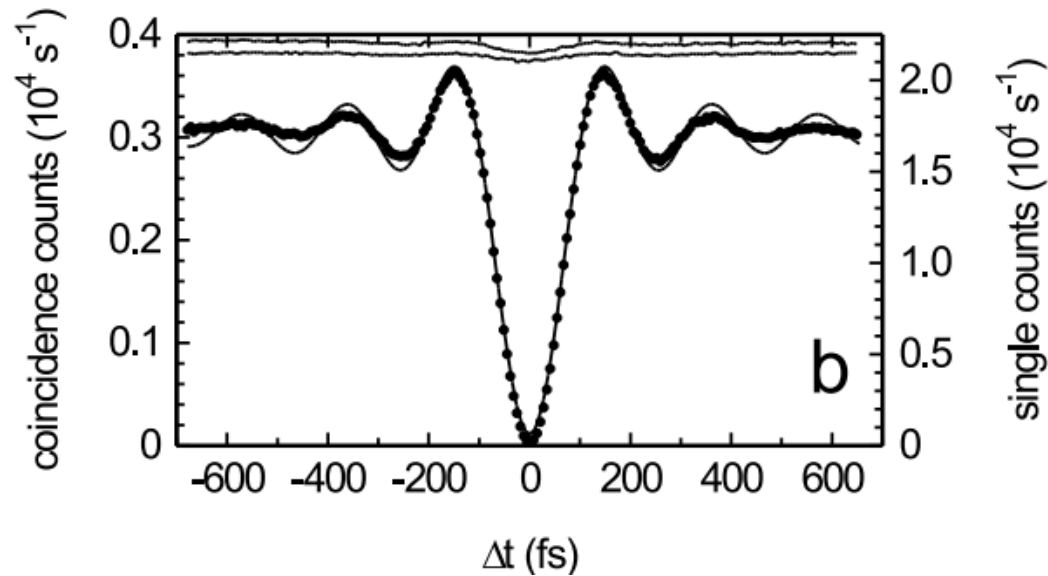
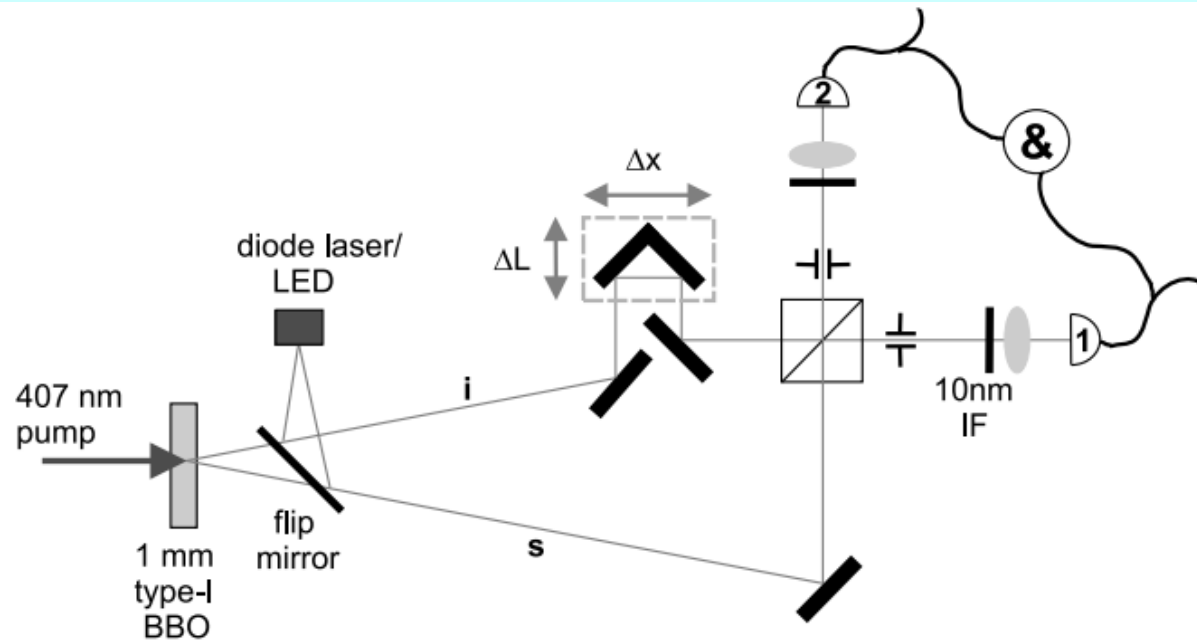


Identical photons:

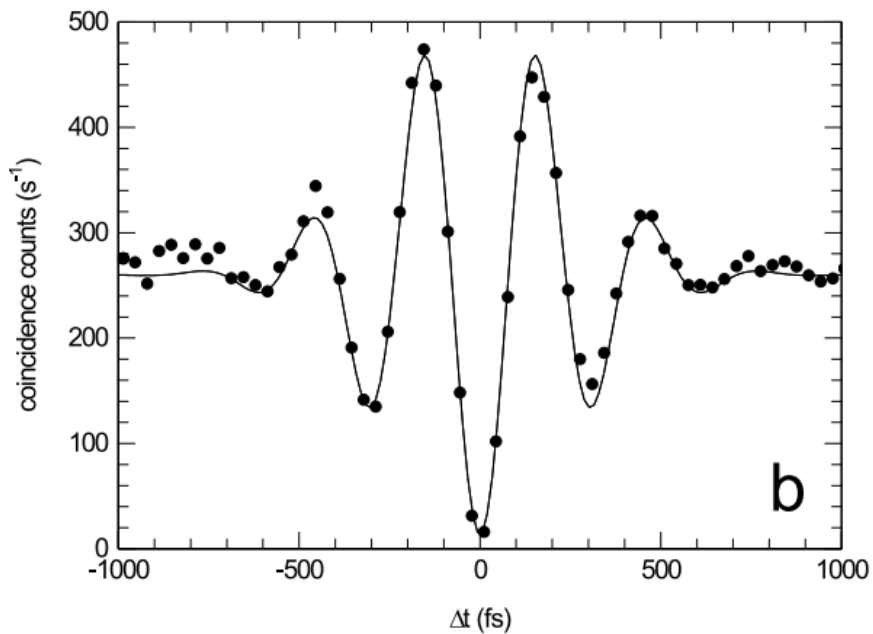
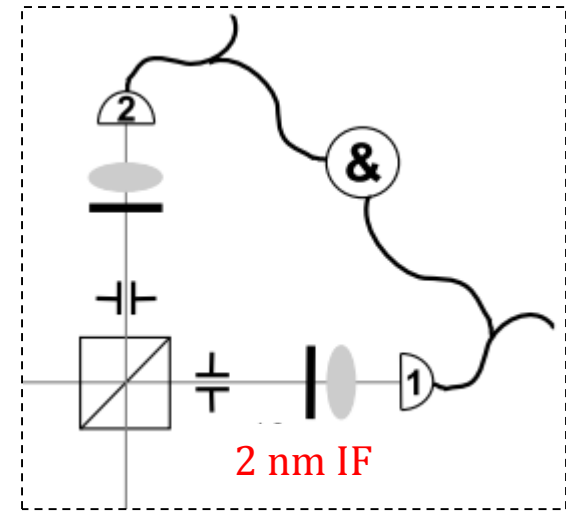
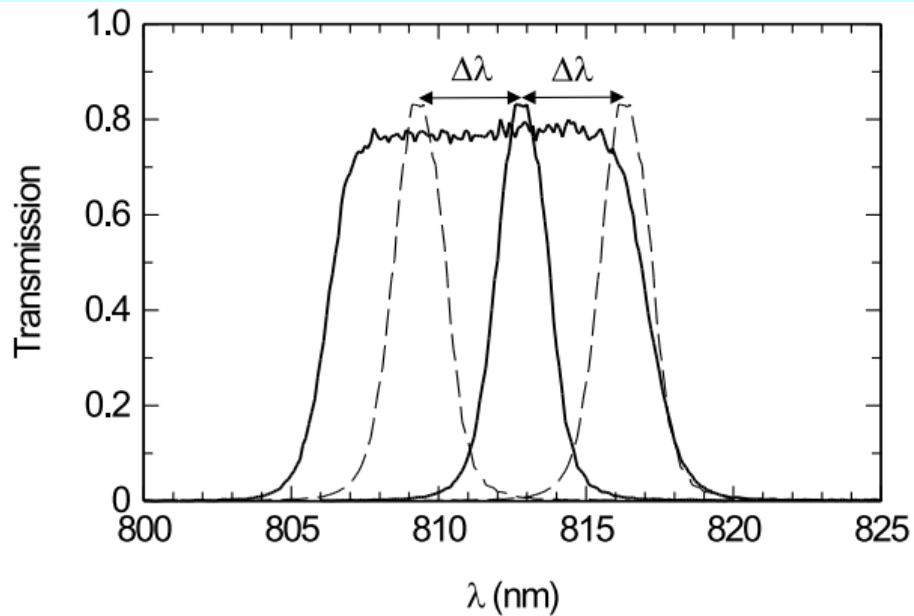
- Same spatial profile (project on single-mode fiber)
- Same time/frequency profile (use pulsed light)

Rarity, Tapster & Loudon (1997)

Our results with HOM interference



Our results with HOM interference



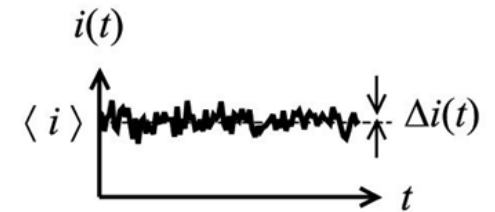
Conclusions: When do you need quantum optics?

- **Single-mode quantum optics**

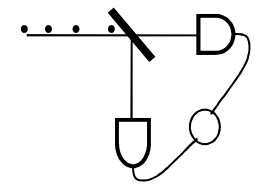
is rarely relevant in free-space optics experiments

- **Multi-mode quantum optics** needs field quantization for:

1. Direct observation of intensity noise (difficult; sensitive to loss)



2. Photon-photon correlations (Hanbury Brown & Twiss)



2b. Interference of two single photons (Hong, Ou & Mandel)

